

## Vibrational Spectra and Normal Coordinate Calculations of $\text{HgNH}_2\text{Cl}$ and $\text{HgNH}_2\text{Br}$ \*

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Infrared spectra of  $\text{HgNH}_2\text{Cl}$  and  $\text{HgNH}_2\text{Br}$  and the Raman spectrum of  $\text{HgNH}_2\text{Cl}$  have been measured in the region from  $4000\text{ cm}^{-1}$  to  $30\text{ cm}^{-1}$ . Effects of changing the  $\text{X}$  ( $\text{X}=\text{Cl}, \text{Br}$ ) ion were observed mainly in N—H stretching bands and in lattice vibrations. Calculations of normal frequencies were carried out using a force field which includes inter-molecular forces. An assignment of the lattice vibrations was proposed and the force constants were also calculated. The force constants of the inter-molecular hydrogen bonds  $\text{N—H}\cdots\text{Cl}$  and  $\text{N—H}\cdots\text{Br}$  were obtained as  $0.212\text{ mdyn/\AA}$  and  $0.142\text{ mdyn/\AA}$ , respectively.

The structure of orthorhombic crystal of mercuric amidohalide  $\text{HgNH}_2\text{X}$  ( $\text{X}=\text{Cl}, \text{Br}$ ) consists of halogen ions and infinite chains of alternating Hg and  $\text{NH}_2$  with linear bonds about Hg and tetrahedral bonds about N.<sup>1–3</sup> The vibrational assignment of the infrared bands observed between  $4000$  to  $400\text{ cm}^{-1}$  has already been given for the infinite chain of  $(\text{HgNH}_2^+)_n$  by Nakagawa *et al.*<sup>4,5</sup> And the force constants associated with the  $(\text{HgNH}_2^+)_n$  chain have also been computed.<sup>4</sup>

However, the lattice vibrations of these compounds have not been studied. It is considered interesting to examine the effects of changing the halogen ion on the vibrational frequencies. The effects are expected to

appear mainly in the lattice vibrations and also in the N—H stretching vibrations because of the possible  $\text{N—H}\cdots\text{X}$  type hydrogen bonds.

We have measured the infrared spectra of  $\text{HgNH}_2\text{X}$  ( $\text{X}=\text{Cl}, \text{Br}$ ) and the Raman spectrum of  $\text{HgNH}_2\text{Cl}$  in the region from  $4000$  to  $30\text{ cm}^{-1}$ . The calculation of the normal frequencies was carried out including  $\text{Cl}^-$  and  $\text{Br}^-$  ions, and hence both the intra-molecular force constants and the inter-molecular ones were obtained. It then becomes possible to discuss quantitatively about the hydrogen bond  $\text{N—H}\cdots\text{X}$  ( $\text{X}=\text{Cl}, \text{Br}$ ) and its effects on the N—H stretching force constant.

### Experimental

Mercuric amidochloride  $\text{HgNH}_2\text{Cl}$  of special grade (Sanko Seiyaku Kogyo Co. Ltd.) was used without further purification. Although mercuric amidobromide  $\text{HgNH}_2\text{Br}$  has two types of modifications, namely orthorhombic and cubic, only the orthorhombic one was prepared by the reaction of  $150\text{ ml}$  of  $14\text{ N}$  aqueous ammonia with  $2\text{ g}$  of mercury bromide  $\text{Hg}_2\text{Br}_2$ .<sup>2,6</sup> The infrared spectrum of the prepared  $\text{HgNH}_2\text{Br}$  coincided with that of the orthorhombic modification reported by Nakagawa *et al.*<sup>5</sup>

The infrared spectra of these substances dispersed in Nujol and in KBr disks were recorded by a Hitachi EPI G3 spectrometer ( $4000$ – $400\text{ cm}^{-1}$ ) and these in Nujol by a Hitachi FIS 3 spectrometer ( $400$ – $30\text{ cm}^{-1}$ ). However, the KBr disk method was not adequate for  $\text{HgNH}_2\text{Cl}$  because it was found that a part of chlorine ions in the sample were exchanged by bromine ions in KBr. The Raman spectrum of  $\text{HgNH}_2\text{Cl}$  crystal powder pressed to a disk was obtained by a laser Raman spectrometer (Japan Electron Optics Laboratory Co. Ltd.). The  $5145\text{ \AA}$  line of argon ion was used as the light source for excitation.

### Observed spectra

The infrared spectra of  $\text{HgNH}_2\text{Cl}$  and  $\text{HgNH}_2\text{Br}$  observed from  $4000$  to  $400\text{ cm}^{-1}$  are shown in Figs. 2 and 3, respectively. The far-infrared spectra from  $400$  to  $30\text{ cm}^{-1}$  are given in Figs. 4 and 5. In Fig. 6, the Raman spectrum of  $\text{HgNH}_2\text{Cl}$  is reproduced.

Vibrations of  $\text{HgNH}_2$ . Nakagawa *et al.*<sup>4,5</sup> have

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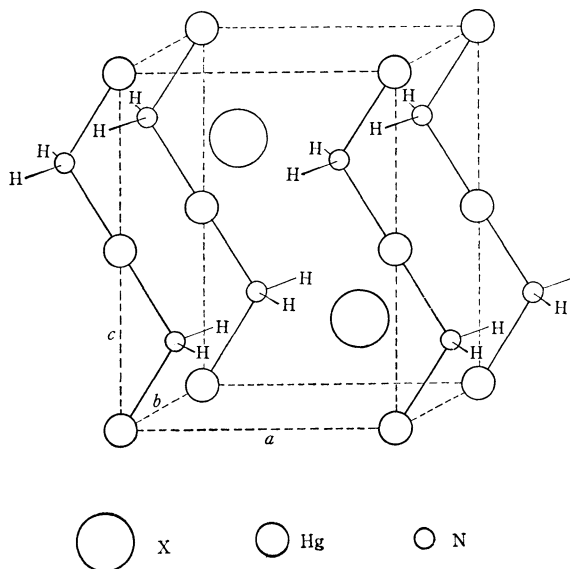


Fig. 1. Crystal structure of  $\text{HgNH}_2\text{X}$  ( $\text{X}=\text{Cl}, \text{Br}$ ).

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1) W. N. Lipscomb, *Acta Crystallogr.*, **4**, 266 (1951).  
 2) L. Nijssen and W. N. Lipscomb, *ibid.*, **5**, 604 (1952).  
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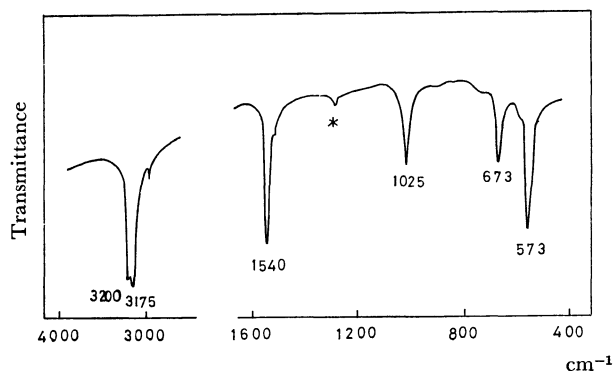


Fig. 2. Infrared spectrum of  $\text{HgNH}_2\text{Cl}$  in nujol mull (nujol bands eliminated). \* is due to  $\text{Hg}(\text{NH}_3)_2\text{Cl}_2$ .<sup>4)</sup>

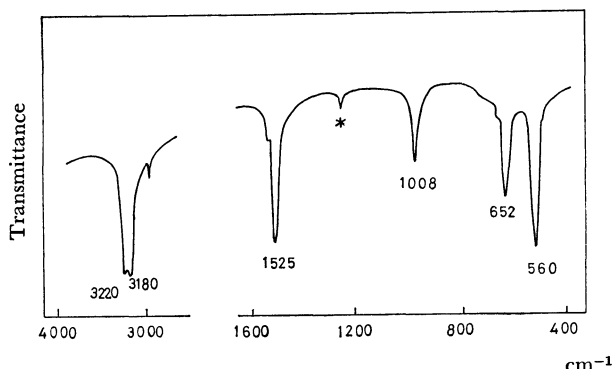


Fig. 3. Infrared spectrum of  $\text{HgNH}_2\text{Br}$  in nujol mull (nujol bands eliminated). \* is due to  $\text{Hg}(\text{NH}_3)_2\text{Br}_2$ .<sup>5)</sup>

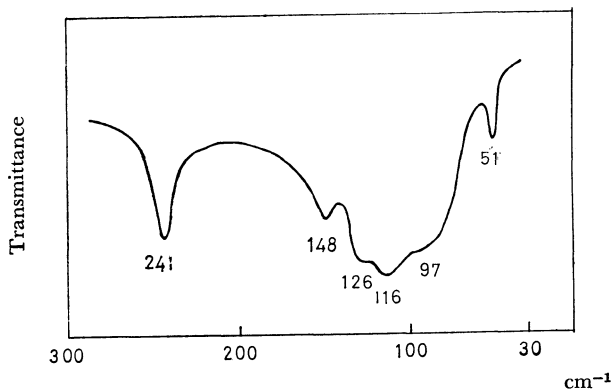


Fig. 4. Far-infrared spectrum of  $\text{HgNH}_2\text{Cl}$  in nujol mull.

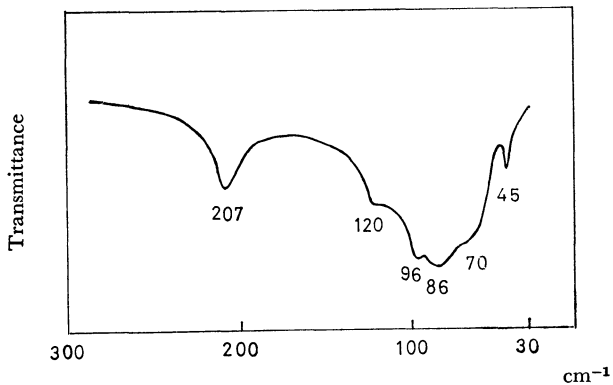


Fig. 5. Far-infrared spectrum of  $\text{HgNH}_2\text{Br}$  in nujol mull.

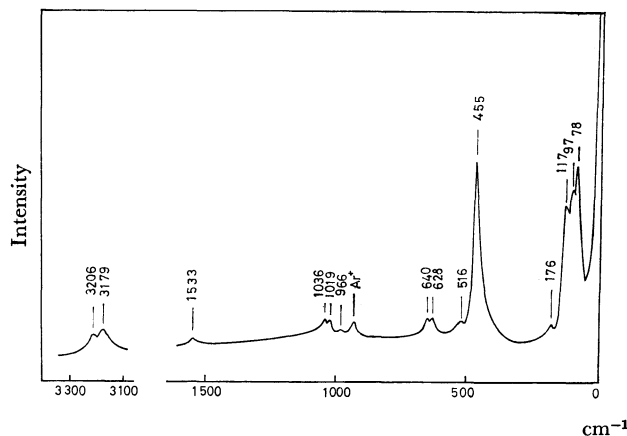


Fig. 6. Raman spectrum of  $\text{HgNH}_2\text{Cl}$  powders (pressed disk).

already measured the infrared spectra of  $\text{HgNH}_2\text{Cl}$  and  $\text{HgNH}_2\text{Br}$  from 4000 to 400  $\text{cm}^{-1}$  and have given the assignment of the infrared active vibrations ( $B_{1u}$ ,  $B_{2u}$  and  $B_{3u}$  symmetry species) of  $\text{HgNH}_2^+$ . The frequencies which we observed and the assignment which we have made are in substantial agreement with theirs, hence the complete list is not reproduced here. Our contribution is to measure Raman spectrum which was not reported previously and to present the assignment of the Raman active vibrations ( $A_g$ ,  $B_{1g}$ ,  $B_{2g}$ , and  $B_{3g}$  symmetry species).

One interesting feature in the infrared spectra, however, should be pointed out. The  $\text{NH}_2$  anti-symmetric stretching vibration is observed at 3200  $\text{cm}^{-1}$  in  $\text{HgNH}_2\text{Cl}$  while it is observed at 3220  $\text{cm}^{-1}$  in  $\text{HgNH}_2\text{Br}$ . The same phenomenon is observed in the  $\text{NH}_2$  symmetric stretching also: 3175  $\text{cm}^{-1}$  for  $\text{HgNH}_2\text{Cl}$  and 3180  $\text{cm}^{-1}$  for  $\text{HgNH}_2\text{Br}$ . These high-frequency shifts in N-H stretching vibrations when  $\text{Cl}^-$  is substituted by  $\text{Br}^-$  are explained by the existence of the  $\text{N-H}\cdots\text{X}$  ( $\text{X}=\text{Cl}, \text{Br}$ ) hydrogen bond. All the other vibrations shift to the side of lower frequencies when  $\text{Cl}^-$  is exchanged by  $\text{Br}^-$ .

In the Raman spectrum of  $\text{HgNH}_2\text{Cl}$ , ten bands are observed above 400  $\text{cm}^{-1}$ . The assignment of these Raman bands presented in Table 2 was made by analogy with that of the infrared bands. However, the observed two doublets of  $\text{NH}_2$  wagging and of  $\text{NH}_2$  rocking vibrations are not explained unambiguously at present. They may be either due to combination bands or to impurities in the sample.

**Chain bending and lattice vibrations.** There has been no work of the far-infrared spectra below 400  $\text{cm}^{-1}$  of these compounds. We have observed the bands due to chain bending vibrations and lattice vibrations of  $\text{HgNH}_2\text{Cl}$  and  $\text{HgNH}_2\text{Br}$  for the first time. They are shown in Figs. 4 and 5, respectively. In the spectrum of  $\text{HgNH}_2\text{Cl}$ , six bands are observed at 241, 148, 126, 116, 97, and 51  $\text{cm}^{-1}$ . The corresponding bands in  $\text{HgNH}_2\text{Br}$  spectrum are 207, 120, 96, 86, 70, and 45  $\text{cm}^{-1}$ . These bands are considered to be due to chain bending vibrations and lattice vibrations or the mixture of these.

In the low frequency region of the Raman spectrum of  $\text{HgNH}_2\text{Cl}$ , four bands are observed at 176, 117, 97,

and 78 cm<sup>-1</sup>. The assignment of these low frequency bands observed in the infrared and Raman spectra are made according to the results of the normal frequency calculations.

### Normal coordinate treatment

The crystal structure of HgNH<sub>2</sub>X (X=Cl, Br) is orthorhombic with two HgNH<sub>2</sub>X units in a Bravais unit cell as shown in Fig. 1. The lattice parameters and the bond-lengths are summarized in Table 1.<sup>1,2)</sup>

TABLE 1. LATTICE PARAMETERS AND BOND LENGTHS<sup>1,2)</sup>

	HgNH <sub>2</sub> Cl	HgNH <sub>2</sub> Br
Lattice parameter	<i>a</i> =5.167 Å <i>b</i> =4.357 <i>c</i> =6.690	<i>a</i> =5.439 Å <i>b</i> =4.487 <i>c</i> =6.761
Bond length		
N—H	1.04 Å	1.04 Å
Hg—N	2.05	2.07
H...H	1.70	1.70
H...Hg	2.59	2.61
Hg...Hg	3.35	3.38
H...X	2.18	2.36
Hg...X	3.21	3.30
X...X	3.83	3.91

The calculation of the normal frequencies<sup>7)</sup> was made assuming a modified Urey-Bradley type force field<sup>8)</sup> and a few coordinates corresponding to inter-molecular forces were also taken into account. These coordinates are given below.

	X=Cl	X=Br
q (H...X)	2.18 Å	2.36 Å
q (Hg...X)	3.21 Å	3.30 Å
h (N-H...X)		

The coordinate q(X...X) was omitted because of the large X...X distance: 3.83 Å for X=Cl and 3.91 Å for X=Br.

The potential function was expressed as

$$U = U_1 + U_2$$

$$U_1 = 1/2 \{ \sum K_{N-H} (\Delta r_{N-H})^2 + \sum K_{Hg-N} (\Delta r_{Hg-N})^2 + \sum F_{H...H} (\Delta q_{H...H})^2 + \sum F_{H...Hg} (\Delta q_{H...Hg})^2 + \sum H_{H-N-H} (\Delta \alpha_{H-N-H})^2 + \sum H_{H-N-Hg} (\Delta \alpha_{H-N-Hg})^2 + \sum H_{Hg-N-Hg} (\Delta \alpha_{Hg-N-Hg})^2 + \sum H_{N-Hg-N}^{\perp} (\Delta \alpha_{N-Hg-N}^{\perp})^2 + \sum H_{N-Hg-N}^{\parallel} (\Delta \alpha_{N-Hg-N}^{\parallel})^2 + \sum Y_{N-Hg-N} (\Delta \tau_{N-Hg-N})^2 \} + \sum \text{angle interaction constant } l_{ij} \text{ term}^9)$$

$$U_2 = 1/2 \{ \sum f_{H...X} (\Delta q_{H...X})^2 + \sum f_{Hg...X} (\Delta q_{Hg...X})^2 + \sum h_{N-H...X} (\Delta \phi_{N-H...X})^2 \}$$

where

$U_1$  is the potential function associated with the (HgNH<sub>2</sub>)<sub>n</sub> chain.

7) T. Shimanouchi, M. Tsuboi, and T. Miyazawa, *J. Chem. Phys.*, **35**, 1597 (1961).

8) T. Shimanouchi, *Pure Appl. Chem.*, **7**, 131 (1963).

9) T. Shimanouchi and I. Suzuki, *J. Mol. Spectrosc.*, **6**, 277 (1961); **8**, 222 (1962).

$U_2$  is that of the lattice vibration.

K, F, H, and Y denote bond-stretching, repulsive, angle-bending and torsional force constants, respectively. f and h are, respectively, stretching and angle-bending force constants associated with the lattice vibrations.

The calculation of the normal frequencies was carried out by the Cartesian coordinate method.<sup>7)</sup> The coordinates, with which the potential function is expressed, are transformed into Cartesian displacement coordinates X. And then the optically active potential energy matrix  $F_{op}$  is set up in terms of the optically active Cartesian displacement coordinates. The secular equation

$$| M^{-1} F_{op} - E \lambda | = 0$$

gives the eigenvalues and eigenvectors which correspond to the frequencies and the vibrational modes respectively, where  $M$  is a diagonal mass matrix and  $E$  is a unit matrix.

### Results and Discussions

**Normal frequency calculation.** Lipscomb *et al.* have reported from X-ray studies that the space groups of both HgNH<sub>2</sub>Cl<sup>1)</sup> and HgNH<sub>2</sub>Br<sup>2)</sup> belong to  $C_{2v}^1$ -P2mm. However, as can be seen in Table 2, there are very few infrared frequencies which correspond to Raman fre-

TABLE 2. OBSERVED AND CALCULATED FREQUENCIES IN cm<sup>-1</sup>

Species	HgNH <sub>2</sub> Cl		HgNH <sub>2</sub> Br		Assignment
	Obsd.	Calcd.	Obsd.	Calcd.	
$A_g$	3179	3176		3185	NH <sub>2</sub> sym. str.
	1533	1539		1535	NH <sub>2</sub> scis.
	455	459		449	Hg-N sym. str. + Hg-N-Hg def.
$B_{1g}$	117	130		77	lattice
	3206	3204		3215	NH <sub>2</sub> anti. str.
	(640 628)	642		615	NH <sub>2</sub> rock.
	176	165		119	lattice
	97	93		66	lattice
$B_{2g}$	(1036 1019)	1018		1004	NH <sub>2</sub> wag.
	516	524		525	Hg-N anti. str.
	78	76		47	lattice
$B_{3g}$	966	960		921	NH <sub>2</sub> twist.
$A_u$	inact.	967	inact.	928	NH <sub>2</sub> twist.
		68		66	N-Hg-N bend. (⊥)
$B_{1u}$	1025	1019	1008	1006	NH <sub>2</sub> wag.
	573	564	560	565	Hg-N anti. str.
	97	93	70	76	N-Hg-N bend. (//)
	51	53	45	44	lattice
$B_{2u}$	3200	3204	3220	3215	NH <sub>2</sub> anti. str.
	673	675	652	649	NH <sub>2</sub> rock.
	241	234	207	216	N-Hg-N bend. (⊥)
	126	135	96	90	lattice
$B_{3u}$	3175	3176	3180	3185	NH <sub>2</sub> sym. str.
	1540	1540	1525	1535	NH <sub>2</sub> scis.
	—	508	—	499	Hg-N sym. str.
	148	141	120	116	lattice
	116	117	86	91	N-Hg-N bend. (//) + Hg-N-Hg def.

quencies. This fact would suggest that the vibrational spectra of  $\text{HgNH}_2\text{Cl}$  (and also  $\text{HgNH}_2\text{Br}$ ) are explained more adequately by a  $D_{2h}$  symmetry than a  $C_{2v}$ . Therefore we have carried out normal frequency calculations assuming a  $D_{2h}$  symmetry instead.

The calculated frequencies listed in Table 2 are in good agreement with the observed values. As the Bravais unit cell contains one  $(\text{HgNH}_2^+)_n$  chain and two X atoms, one rotational and six translational lattice vibrations are expected to appear. The assignment of lattice vibrations and also of chain bending vibrations are given in the same table. The  $L_x$  matrices of lattice vibrational modes and chain bending modes are shown schematically in Figs. 7, 8, and 9.  $L_x$  is the transfor-

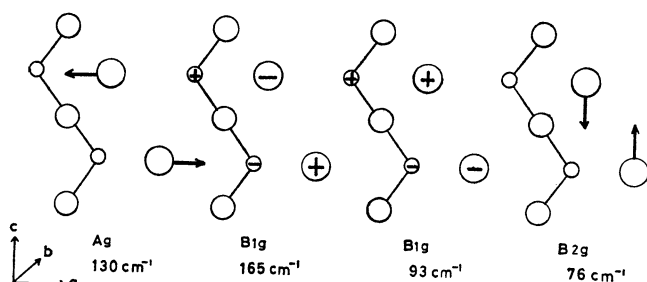


Fig. 7. Displacements of atoms for the Raman active lattice vibrational modes, with calculated frequencies of  $\text{HgNH}_2\text{Cl}$ .

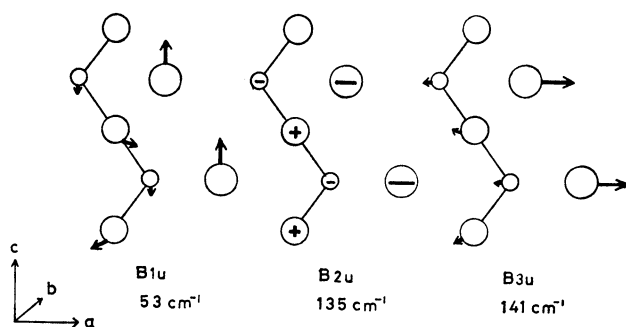


Fig. 8. Displacements of atoms for the infrared active lattice vibrational modes, with calculated frequencies of  $\text{HgNH}_2\text{Cl}$ .

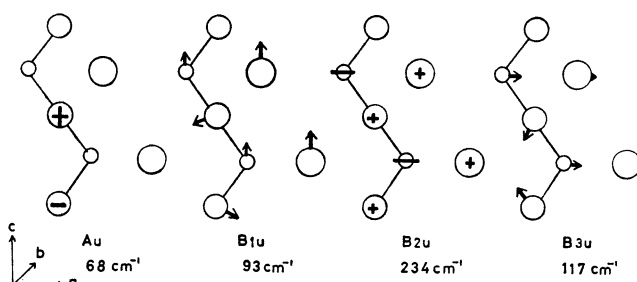


Fig. 9. Displacements of atoms for the chain bending modes, with calculated frequencies of  $\text{HgNH}_2\text{Cl}$ .

mation matrix between the Cartesian coordinates  $X$  and the normal coordinates  $Q$ . These figures demonstrate that the low frequency bands cannot be distinctly classified into either lattice vibrations or chain bending vibrations; rather, they are more or less coupled with each other. The low-frequency shifts in the chain bending vibrations, when  $\text{Cl}^-$  is substituted by  $\text{Br}^-$ , are well understood from this viewpoint. The chain bending vibrations in the  $A_u$  species of  $\text{HgNH}_2\text{Cl}$  and

$\text{HgNH}_2\text{Br}$ , however, are both calculated to be almost of the same value:  $68\text{ cm}^{-1}$  for  $\text{HgNH}_2\text{Cl}$  and  $66\text{ cm}^{-1}$  for  $\text{HgNH}_2\text{Br}$ . As can be seen in Table 2 and in Fig. 9, this mode is a pure chain bending vibration so that the effect of changing the halogen ion is negligible. Low-frequency shifts in the lattice vibrations when  $\text{Cl}^-$  is exchanged by  $\text{Br}^-$  are explained mainly both by the difference of atomic weight of halogen and by the difference of inter-molecular force constants which will be discussed below.

TABLE 3. FORCE CONSTANTS IN  $\text{mdyn}/\text{\AA}$

	$\text{HgNH}_2\text{Cl}$	$\text{HgNH}_2\text{Br}$
$K(\text{N-H})$	5.276	5.371
$K(\text{Hg-N})$	1.775	1.775
$F(\text{H}\cdots\text{H})$	0.100	0.100
$F(\text{H}\cdots\text{Hg})$	0.160	0.160
$f(\text{H}\cdots\text{X})$	0.212	0.142
$f(\text{Hg}\cdots\text{X})$	0.111	0.101
$H(\text{H-N-H})$	0.555	0.555
$H(\text{H-N-Hg})$	0.105	0.105
$H(\text{Hg-N-Hg})$	0.081	0.081
$h(\text{N-H}\cdots\text{X})$	0.016	0.011
$H^\perp(\text{N-Hg-N})$	0.083	0.083
$H^{\parallel}(\text{N-Hg-N})$	0.077	0.077
$Y(\text{N-Hg-N})$	0.030	0.030 <sup>a)</sup>
$l(\text{Hg, H, H})$	-0.003	0.002 <sup>a)</sup>
$l(\text{H, H, Hg})$	0.009	0.009 <sup>a)</sup>
$l(\text{H, Hg, Hg})$	-0.035	-0.024 <sup>a)</sup>

a)  $\text{mdyn}\cdot\text{\AA}$

$\text{N-H}\cdots\text{X}$  ( $\text{X}=\text{Cl, Br}$ ) hydrogen bond. The values of the force constants obtained by the least squares method are shown in Table 3. In the course of the calculation, all the force constants except for  $K(\text{N-H})$ ,  $f(\text{H}\cdots\text{X})$ ,  $f(\text{Hg}\cdots\text{X})$ ,  $h(\text{N-H}\cdots\text{X})$  and four kinds of angle interaction constant  $l^0$  were set to be equal for  $\text{HgNH}_2\text{Cl}$  and  $\text{HgNH}_2\text{Br}$  in spite of the slight difference of the bond lengths summarized in Table 1. Nakagawa *et al.* determined  $F(\text{H}\cdots\text{H})$  to be  $0.06\text{ mdyn}/\text{\AA}$ <sup>4)</sup>, however, in this calculation it was fixed to  $0.10\text{ mdyn}/\text{\AA}$  according to Harada and Shimanouchi.<sup>10)</sup>

The force constants  $K(\text{N-H})$  are calculated as  $5.276\text{ mdyn}/\text{\AA}$  for  $\text{HgNH}_2\text{Cl}$  and  $5.371\text{ mdyn}/\text{\AA}$  for  $\text{HgNH}_2\text{Br}$ . High-frequency shifts in the N-H stretching vibrations when  $\text{Cl}^-$  is substituted by  $\text{Br}^-$  are explained by the fact that  $K(\text{N-H})$  of  $\text{HgNH}_2\text{Br}$  is larger than that of  $\text{HgNH}_2\text{Cl}$ . This difference in  $K(\text{N-H})$  between  $\text{HgNH}_2\text{Cl}$  and  $\text{HgNH}_2\text{Br}$  may suggest that the hydrogen bond of  $\text{H}\cdots\text{Cl}$  is stronger than that of  $\text{H}\cdots\text{Br}$ . And actually, the calculated values of the force constant  $f(\text{H}\cdots\text{Cl})$  and  $h(\text{N-H}\cdots\text{Cl})$  become larger than those of  $f(\text{H}\cdots\text{Br})$  and  $h(\text{N-H}\cdots\text{Br})$ , respectively, as shown in Table 3. Therefore, low-frequency shifts in the lattice vibrations when  $\text{Cl}^-$  is exchanged by  $\text{Br}^-$  are understandable. All the other force constants obtained are considered reasonable.

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